

**Delaying future
sea-level rise by
storing water on
Antarctica**

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Delaying future sea-level rise by storing water on Antarctica

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The Antarctic ice sheet is situated on the coldest continent on Earth with most of its surface temperatures far below the freezing point of ocean water throughout the year. The water volume that is equivalent to one meter of global sea-level rise would elevate the Antarctic ice sheet by ~ 25 m if distributed uniformly. The currently observed $\sim 3 \text{ mm yr}^{-1}$ of global average sea-level rise due to thermal expansion, additional water added from glaciers and ice sheets, and changes in land water storage corresponds to about $10^{12} \text{ m}^3 \text{ yr}^{-1}$ of ocean-water volume. Solely in terms of throughput mitigating a sea-level rise of 3 mm yr^{-1} would require 90 of the largest pump stations currently under construction in New Orleans each assumed to pump $\sim 360 \text{ m}^3 \text{ s}^{-1}$ which corresponds to $\sim 11 \times 10^9 \text{ m}^3 \text{ yr}^{-1}$ (Alyeska Pipeline Service Company, 2013). The height of the ice sheet of about 4000 m means that it would require a constant power of 1275 GW to elevate the potential energy of the associated ocean water. This is equivalent to $\sim 7\%$ of the global primary energy supply of the year 2012 (International Energy Agency, 2014). The power required for the actual pumping may even be higher and reach 2300 GW under optimistic assumptions (see Supplement). It will have to be generated by renewable resources to avoid the additional climate change and sea-level rise associated with fossil fuels. The Antarctic continent is windy enough to support such pumping through wind energy, with around 16.7 TW available in a 200 km wide band along the coast of East Antarctica (Archer and Jacobson, 2005) (see Supplement). Around 8% of that energy would need to be extracted to compensate the potential energy increase of the pumped water alone, which is equivalent to 850 000 wind-energy plants of 1.5 MW running on full capacity. Although the approach may be the only way to protect entire coast lines it will not be feasible without major technical innovations solving the fundamental energy problem. In the following we explore the option from an ice dynamical point of view.

Antarctica's currently observed ice loss occurs near the coast (Shepherd et al., 2012) while the surface in its interior is moving at a speed of less than 0.1 m yr^{-1} (Rignot et al., 2011). Because the ice is continually moving, ocean water put on the ice sheet will only

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To represent the large-scale dynamics reasonably well, we use a 12 km horizontal resolution for the ice sheet simulations. Our hybrid shallow approximation ensures stress transmission across the grounding line and a smooth transition between regimes of fast flowing, sliding ice and slowly deforming, bedrock-frozen ice. The grounding line can freely evolve also under lower resolution due to a local interpolation of the grounding-line position that affects the basal friction and a new driving stress scheme at the grounding line. The interpolation leads to a reversible grounding-line dynamics that is consistent with Full-Stokes simulations in high resolution (Feldmann et al., 2014). Although the model is capable of simulating the coastal dynamics of the ice sheet within limitations, it is important to note that the results obtained here are predominantly dependent on the ice flow representation in the interior of the ice sheet for that large-scale continental ice-sheet models like PISM and others (Bindshadler et al., 2013; Calov et al., 2010; Greve et al., 2011; Huybrechts and Wolde, 1999; Pollard and Deconto, 2009; Swingedouw et al., 2008) have been designed.

In the standard simulation we do not alter the surface air temperature during the mass addition. To estimate the effect of surface warming due to the latent heat release of the sea water we conduct a second set of simulations that keeps the surface temperature at the freezing point of sea water during the surface mass addition.

3 Results

The ratio of the volume added during the first 100 years and the volume that is lost again after 1000 years depends strongly on the distance from the coast (Figs. 2 and 3). Consistent with earlier studies (Huybrechts and Wolde, 1999; Winkelmann et al., 2012) an ice volume equivalent to 10–15 % of the added ice is already lost at the end of the forcing period when the ice is added at a distance of 200 km from the coastline, while the sea-level contribution is strongly delayed at a distance above 500 km from the coast (Fig. 2). In order to minimize the return flow of the ice into the ocean, the specific positioning of the ice addition could be varied spatially making use of slow moving ice

regions. Here we apply a simplified spatial distribution in order to demonstrate the main process and allow for a conceptual analysis of the simulations.

The time after which a certain equivalent of the added ice has been discharged into the ocean, is much shorter than would be expected from a mere advection of the added ice mass with the surface velocities of the ice sheet (Fig. S2 in the Supplement). That is because the ice thickness anomaly creates an imbalance between the driving stress and the viscous ice flow. As a consequence the ice transport occurs in waves from the strip of perturbation to the coast (Winkelmann et al., 2012) (Fig. 1) with the ice being discharged to the ocean not being the same ice that was added to the ice sheet earlier. Even though the ice wave also travels partially inland, it is possible that more ice is transported out of the continent than was initially added. This ice-loss exceedance, however, occurs only several millennia after the perturbation (Fig. S2). Whether it is directly related to the perturbation or a manifestation of a localized multistability of the ice dynamics is difficult to identify, because the differences between the initial and final ice topography are within the uncertainty range of the model performance.

The freezing of ocean water that is pumped onto the ice surface will release latent heat that will heat up both the atmosphere and the upper ice layer. The maximum injection of latent heat occurs for the 10 mm yr^{-1} sea-level-mitigation scenario and the 800 km band, which has the smallest area. The corresponding addition of 3.2 m yr^{-1} liquid sea water is equivalent to a latent heat injection of about 35 W m^{-2} . In order to estimate the effect of the surface warming on the ice dynamics, we conduct a second set of simulations in which the surface temperature is held at the freezing point of sea water of -1.9°C during the forcing. This imitates the situation in which the surface of the ice remains in a mixed state of ice and water. As Antarctica's inland-surface temperatures are far below zero, this constitutes a strong warming signal that diffuses down into the ice body and causes ice to soften and flow faster.

A warming from -20 to -1.9°C would increase the long-wave radiative loss to the atmosphere by 70 W m^{-2} according to the Stefan–Boltzmann law assuming an emissivity of 0.95. If open water areas are sustained on the ice sheet, a sensible-

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heat-dominated loss can remove heat at a rate of 100 W m^{-2} or more as observed in ocean polynyas (Launiainen and Vihma, 1994). The maximum rate of latent heat injection of 35 W m^{-2} is much smaller than the potential of the atmosphere to remove the heat. Thus keeping surface ice temperatures at freezing point underestimates the atmospheric heat loss so that the simulations provide an upper bound for the induced warming of ice. For perturbation areas with large distance from the coast, the discharge rate of the ice sheet (Fig. 4, thin lines) is nearly identical to the response without warming (Fig. 4, thick lines) on the millennial time scale considered. Within the first millennium, the latent heat release of freezing sea water only alters the discharge when placed near the coast. After 2000 years, the additional warming can induce a discharge exceeding 100 % of the added ice in the 200 km simulations (Fig. S1).

It is possible that ice dynamic effects which are not included in these simulations (such as ice fractures or basal sliding conditions) alter the results quantitatively, but the shallow ice approximation which dominates the ice dynamics in the model in the interior of Antarctica has been shown to represent the interior ice sheet flow on multi-centennial and longer time scales (Greve and Blatter, 2009).

4 Discussion

All scenarios considered here assume that the only perturbation of the ice sheet is the addition of ice mass in bands of the interior of East Antarctica. At the same time Antarctica's coastal regions are out of balance in a number of regions predominantly in West but also in East Antarctica. In this study it is assumed that the addition of ice in the interior will not interfere with the imbalance at the coast. This might be an over-simplification, but currently available modeling studies (Favier et al., 2014; Joughin et al., 2014; Mengel and Levermann, 2014) indicate that perturbations near the coast will not reach as far inland over time periods of several centuries. A possible interference between the interior of the ice sheet and its coastal regions, however, needs further investigation, possibly with higher-resolution regional ice sheet models.

coastal ecosystems. It has to be investigated how the water extraction will influence the small- and large-scale ocean circulation. The ice-rheological changes that are induced by the addition of salt water have to be investigated together with potential effects on the basal conditions of the ice.

5 The heat released from freezing and the pumping process itself is in the order of 10 TW (latent heat) + 1TW (heat released from pumping). This corresponds to about 10 % of the maximum increase in latent heat transport in high northern latitudes under an SRES A1B transient simulation (Held and Soden, 2006). The latent heat release is considered as a main contribution to the Arctic amplification of global warming. Put
10 into this perspective the pumping-induced energy over Antarctica is not negligible but significantly smaller than the warming induced latent heat released in northern high latitudes. Potential consequences for the atmospheric and oceanic circulation have to be further explored.

15 When the pumping is stopped the additional discharge from Antarctica will increase the rate of sea-level rise even beyond the warming-induced rate. In this sense the presented approach means raising a loan on Antarctica that future generations will have to pay back. In all simulations considered here pumping is stopped after 100 years, i.e. it is investigated as an option to delay part of the sea-level rise we are already committed to but not as a permanent measure that may induce further
20 responses of the ice sheet not captured here.

If at all feasible, the considered scenarios do not at all represent an alternative to the mitigation of carbon emissions because the method does not address any other climate-change impact than sea-level rise. Furthermore unmitigated emission might induce a sea-level rise of 10 mm yr^{-1} and beyond, which increases the impacts on
25 Antarctica and the burden for future generations when mitigated by pumping of ocean waters. Thirdly, after the pumping is stopped, the sea level will accelerate quickly towards the rate that corresponds to the warming level plus that induced by the addition of ice onto Antarctica.

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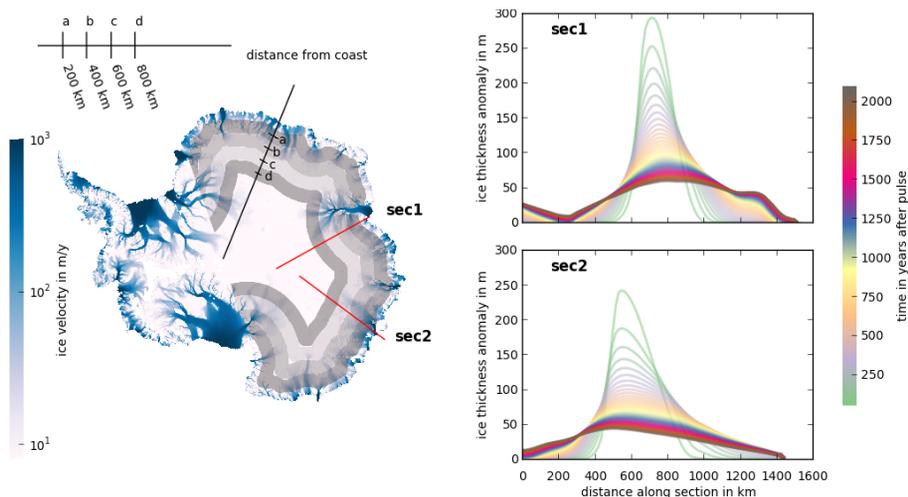


Figure 1. Bands of ice mass addition on East Antarctica and ice thickness relaxation for the 800 km band. Left panel: surface velocities of the ice flow of the Antarctic ice sheet (blue shading). Grey strips indicate where ice mass was added to East Antarctica in order to delay future sea-level rise in the different simulations. The ice was added in strips of 200 km width for 100 years. The right panels show the ice thickness relaxation after the end of the mass addition to the 800 km band in time steps of 50 years for two representative sections (left panel, red lines) as anomaly to the equilibrium simulation.

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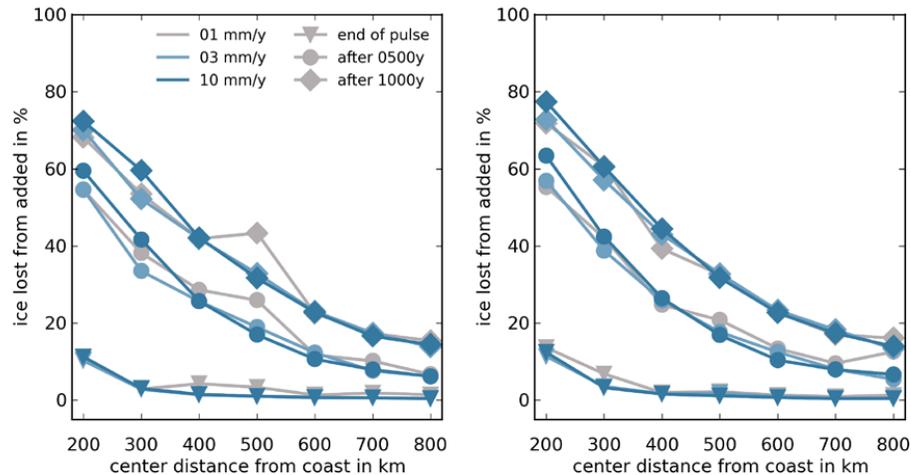


Figure 2. Additional ice discharge compared to mass addition for different distances from the coast. The fraction of the added ice that is lost again to the ocean as a function of distance from the coast at which the additional ice was placed for the simulations without (left panel) and with (right panel) surface warming. Colours indicate the magnitude of the mass addition equivalent to 1, 3 and 10 mm yr^{-1} sea level rise mitigation. Markers correspond to the different relaxation times: end of pulse, 500 and 1000 years after the pulse of mass addition ended.

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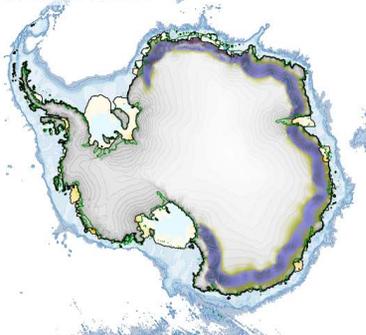
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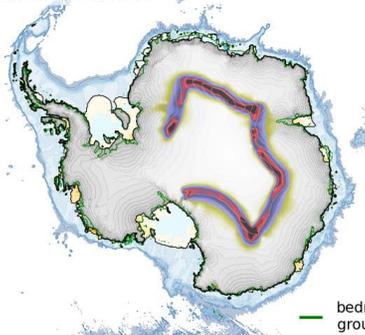
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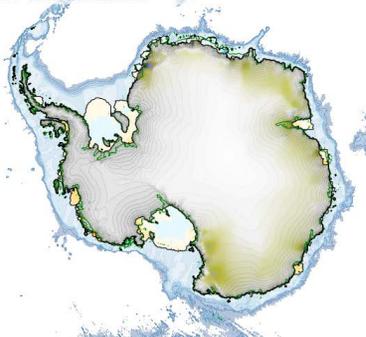
end of 100y pulse
200 km distance



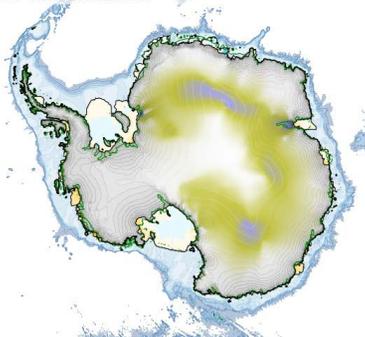
end of 100y pulse
800 km distance



1000y after pulse end
200 km distance



1000y after pulse end
800 km distance



— bedmap2
grounding line
— modeled
grounding line

300
270
240
210
180
150
120
90
60
30
0
ice gain from equilibrium [m]

Figure 3. Difference in ice thickness compared to the initial state. Ice thickness gain at the end of the 100-year-long mass addition (upper panels) and 1000 years after the forcing ended (lower panels). The close-to-coast simulation (left panels) has lost most of the added ice to the ocean after 1000 years while there is a broad ice gain in the 800 km simulation (right panels). Figures are shown for the strongest scenario of 10 mm yr^{-1} of sea-level mitigation and without accounting for latent-heat release.

